## Lecture 2

## PROPERTIES OF FLUID

## Learning Objectives

Upon completion of this chapter, the student should be able to:

- Define three states of matter: Solid, liquid and gas.
- Define mass density, specific weight and weight density.
- Understand the meaning of the term pressure.
- Explain the difference between gauge and absolute pressures.
- Understand the difference between kinematic and absolute viscosities.
- Differentiate between the ideal and real fluids.
- Define bulk modulus
- Explain the dependence of viscosity on temperature.


### 1.1 Introduction

Fluids, both liquids and gases, are characterized by their continuous deformation when a shear force, however small, is applied. Liquids and gases may be distinguished by their relative incompressibilities and the fact that liquid may have a free surface while a gas expands to fill its confining container. Because the liquid is the medium of transmission of power in hydraulic system, knowledge of its characteristics is essential. The purpose of this chapter is to define certain fundamental properties of fluids which will be useful to apply the basic principles of fluid mechanics to the solution of practical problems on fluid power.

### 1.2 Solids and Fluids (Liquids and Gases)

### 1.2.1 Distinction between a Solid and a Fluid

The molecules of a solid are usually closer to each other than those of a fluid. The attractive forces between the molecules of a solid are so large that a solid tends to retain its shape. This is not the case with a fluid where the attractive forces between the molecules are smaller. An ideal elastic solid deforms under load, and once the load is removed, it returns to its original state.

### 1.2.2 Distinction between a Gas and a Liquid

A fluid may be either a gas or a liquid. The molecules of a gas are much farther apart than those of a liquid. Hence, a gas is very compressible, and when all external pressure is removed, it tends to expand indefinitely. A liquid is relatively incompressible, and if all pressure, except its own vapor pressure, is removed, the cohesion between molecules holds them together, so that the liquid does not expand indefinitely. Therefore, a liquid may have a free surface, that is, a surface from which all pressure is removed, except its own vapor pressure.

### 1.3 Density, Specific Weight, Specific Volume and Specific Gravity

### 1.3.1 Density

Density $(\rho)$ is defined as mass per unit volume, that is,

$$
\begin{equation*}
\rho=\frac{\text { Mass }}{\text { Volume }}=\frac{m}{V}\left(\frac{\mathrm{~kg}}{\mathrm{~m}^{3}}\right) \tag{1.1}
\end{equation*}
$$

The density characteristics of typical hydraulic fluids are given in Table 1.1.
Table 1.1 Density characteristics of hydraulic fluids

| Fluid | Density <br> $\left(\mathbf{k g} / \mathbf{m}^{3}\right)$ |
| :--- | :--- |
| Shell Tellus ISO 32 mineral oil | 875 |
| Shell HFB 60\%oil,40\%oil | 933 |
| Shell HFC 60\%glycol,40\%water | 1084 |
| Shell HFD phosphate ester | 1125 |
| Shell Naturelle HFE 32 | 918 |

It is well known that density increases with pressure and decreases with temperature. Figure 1.1 shows such a characteristic for an ISO 32 mineral oil.


Figure 1.1 Characteristics of ISO 32 mineral oil.

### 1.3.2Specific Weight

Specific weight $(\gamma)$ is defined as weight per unit volume, that is,

$$
\begin{equation*}
\gamma=\frac{\text { Weight }}{\text { Volume }}\left(\frac{\mathrm{N}}{\mathrm{~m}^{3}}\right) \tag{1.2}
\end{equation*}
$$

Density and specific weight are related by

$$
\rho=\frac{\gamma}{g} \Rightarrow \gamma=\rho g
$$

where $g$ is acceleration due to gravity. Now

$$
\text { Dimension of } \rho=\frac{\text { Dimension of } \gamma}{\text { Dimension of } g}=\frac{\mathrm{N} / \mathrm{m}^{3}}{\mathrm{~m} / \mathrm{s}^{2}}=\frac{\left(\mathrm{kgm} / \mathrm{s}^{2}\right) / \mathrm{m}^{3}}{\mathrm{~m} / \mathrm{s}^{2}}=\frac{\mathrm{kg}}{\mathrm{~m}^{3}}
$$

Note that the density $(\rho)$ is absolute, since it depends on mass, which is independent of location. The specific weight ( $\gamma$ ), on the other hand, is not absolute, since it depends on the value of gravitational acceleration (g), which varies with location, primarily latitude and elevation above mean sea level. The densities and specific weights of fluids vary with temperature. Table 1.2 gives specific weight of common fluids. The specific weight of a liquid varies only slightly with pressure, depending on the bulk modulus of the liquid; it also depends on temperature and the variation may be considerable. Since the specific weight $(\gamma)$ is equal to $\rho g$, the specific weight of a fluid depends on the local value of the acceleration due to gravity in addition to the variation in temperature and pressure.

Table 1.2 Specific weight $\left(\mu \mathrm{nkN} / \mathrm{m}^{3}\right)$ of common liquids at $20^{\circ} \mathrm{C}, 1013$ millibar abs with $g=$ $9.81 \mathrm{~m} / \mathrm{s}^{2}$

| Carbon tetrachloride | 15.6 |
| :--- | :--- |
| Ethyl alcohol | 7.76 |
| Gasoline | 6.6 |
| Glycerin | 12.3 |
| Kerosene | 7.9 |
| Motor oil | 8.5 |
| Seawater | 10.03 |
| Water | 9.79 |

### 1.3.3Specific Volume

Specific volume (SV) is the volume occupied by a unit mass of fluid. We commonly apply it to gases and usually express it in $\mathrm{m}^{3} / \mathrm{kg}$. Specific volume is the reciprocal of density. Thus,

$$
\begin{equation*}
\mathrm{SV}=\frac{1}{\rho} \tag{1.3}
\end{equation*}
$$

### 1.3.4Specific Gravity

Specific gravity (SG) of a given fluid is defined as the specific weight of the fluid divided by the specific weight of water, that is

$$
\begin{gather*}
\mathrm{SG}_{\text {oil }}=\frac{\gamma_{\text {oil }}}{\gamma_{\text {water }}} \\
\Rightarrow \mathrm{SG}_{\text {oil }}=\frac{\rho_{\text {oil }}}{\rho_{\text {water at standard temperature }}} \tag{1.4}
\end{gather*}
$$

SG of a liquid is a dimensionless ratio. Physicists use $4^{\circ} \mathrm{C}$ as the standard temperature, but engineers often use $15.56^{\circ} \mathrm{C}$. In the metric system, the density of water at $4^{\circ} \mathrm{C}$ is $1.00 \mathrm{~g} / \mathrm{cm}^{3}$, equivalent to $1000 \mathrm{~kg} / \mathrm{m}^{3}$, and hence the SG (which is dimensionless) of a liquid has the same numerical value as its density expressed in $\mathrm{g} / \mathrm{mL}$ or $\mathrm{mg} / \mathrm{m}^{3}$.

The SG of a gas is the ratio of its density to that of either hydrogen or air at some specified temperature and pressure, but there is no general agreement on these standards, and so we must explicitly state them in any given case. Since the density of a fluid varies with temperature, we must determine and specify specific gravities at a particular temperature.

## Example 1.1

Air at $20^{\circ} \mathrm{C}$ and atmospheric pressure has a density of $1.23 \mathrm{~kg} / \mathrm{m}^{3}$. Find its specific gravity. What is the ratio of the specific gravity of water to the specific gravity of air at $20^{\circ} \mathrm{C}$ and atmospheric pressure? What is the significance of the ratio?
Solution: Given $T=20^{\circ} \mathrm{C}$, air density $\left(\rho_{\text {air }}\right)=1.23 \mathrm{~kg} / \mathrm{m}^{3}$, water density
$\left(\rho_{\text {water }}\right)=1000 \mathrm{~kg} / \mathrm{m}^{3}$. Specific gravity of air is given by

$$
\frac{\rho_{\text {air }}}{\rho_{\text {water }}}=1.23 \times 10^{-3}
$$

We know that specific gravity of water $=1$. So

$$
\frac{\text { Specific gravity of air }}{\text { Specific gravity of water }}=\frac{1}{1.23 \times 10^{-3}}=813
$$

Therefore, water is 813 times heavier than air.

## Example1.2

The specific weight of oil mixture at ordinary pressure and temperature is $9.81 \mathrm{kN} / \mathrm{m}^{3}$. The specific gravity of mercury is 13.56 . Compute the density of oil mixture and the specific weight and density of mercury.

Solution: Given specific weight of oil mixture $\gamma_{\text {oil }}=9.81$. The density of oil is

$$
\rho_{\text {oil }}=\frac{\gamma_{\text {oil }}}{g}=\frac{9.81}{9.81}=1.00 \mathrm{mg} / \mathrm{m}^{3}=1.00 \mathrm{~g} / \mathrm{mL}
$$

The specific weight and density of mercury are, respectively,

$$
\begin{gathered}
\gamma_{\text {mercury }}=S_{\text {mercury }} \gamma_{\text {oil }}=13.56(9.81)=133.0 \mathrm{kN} / \mathrm{m}^{3} \\
\rho_{\text {mercury }}=S_{\text {mercury }} \rho_{\text {oil }}=13.56(1.00)=13.56 \mathrm{kN} / \mathrm{m}^{3}
\end{gathered}
$$

## Example 1.3

A cylinder container has a diameter of 0.5 m and a height of 1 m . If it is to be filled with a liquid having a specific weight of $2000 \mathrm{~N} / \mathrm{m}^{3}$, how many kg of this liquid must be added?

Solution: Given diameter $(d)=0.5 \mathrm{~m}$, height $(h)=1 \mathrm{~m}$. The volume is given by

$$
\text { Volume }(V)=\frac{\pi d^{2} h}{4}=0.19635 \mathrm{~m}^{3}
$$

Also it is given that specific weight $(\gamma)$ is $2000 \mathrm{~N} / \mathrm{m}^{3}$. Now

$$
\text { Weight }=V \times \gamma=0.19635 \times 2000=392.7 \mathrm{~N}
$$

The mass is given by

$$
\text { Mass }=\frac{W}{g}=40 \mathrm{~kg}
$$

## Example 1.4

One liter of SAE30 oil weighs 8.70 N . Calculate its specific weight, density and specific gravity.

Solution: Given volume $(V)=1$ liter $=10^{-3} \mathrm{~m}^{3}$, weight $(W)=8.7 \mathrm{~N}$. The specific weight is

$$
\gamma=\frac{W}{V}=\frac{8.7}{10^{-3}}=8700 \mathrm{~N} / \mathrm{m}^{3}
$$

and mass density is

$$
\rho_{\text {oil }}=\frac{\gamma}{g}=887.755 \mathrm{~kg} / \mathrm{m}^{3}
$$

Therefore,

$$
\mathrm{SG}_{\text {oil }}=\frac{\rho_{\text {oil }}}{\rho_{\text {water a t standard temperature }}}=\frac{8878.75}{1000}=0.889
$$

### 1.4 Pressure

Pressure is defined as force per unit area. It is the amount of force acting over a unit area, that is

$$
\begin{equation*}
P=\frac{\text { Force }}{\text { Area }}=\frac{F}{A} \tag{1.5}
\end{equation*}
$$

The pressure developed at the bottom of a column of any liquid is called hydrostatic pressure and is given by

$$
\begin{equation*}
P=\rho g h \tag{1.6}
\end{equation*}
$$

where $P$ is the hydrostatic pressure or the pressure at the bottom of liquid column in Pascal or $\mathrm{N} / \mathrm{m}^{2}, \rho$ is the density of liquid in $\mathrm{kg} / \mathrm{m}^{3}, g$ is the acceleration due to gravity in $\mathrm{m} / \mathrm{s}^{2}$ and $h$ is the level of the column of liquid in meters.

Since

$$
P=\rho g h=\frac{m g h}{V}=\frac{W h}{V}
$$

and

$$
P=\gamma h
$$

where $\gamma$ is the specific weight of liquid in $\mathrm{N} / \mathrm{m}^{3}$, we have $P\left(\mathrm{~N} / \mathrm{m}^{2}\right)=\gamma\left(\mathrm{N} / \mathrm{m}^{3}\right) \times h(\mathrm{~m})$.

### 1.4.1 Pressure at the Bottom of a Column of Liquid

Let us now refer to Fig. 1.2, which shows the pressure head developed at the bottom due to the column of liquid. Let $h$ be the height of the liquid column and $W$ be the weight of the liquid. Let the liquid have a specific weight $\gamma$ and volume $V$ :

$$
\begin{align*}
P= & \frac{F}{A}=\frac{W}{A}=\frac{W}{A} \times \frac{V}{V}=\frac{W}{V} \times \frac{V}{A} \\
& =\frac{(\gamma \times V)}{A}=\frac{(\gamma \times A \times H)}{A} \\
\hline P & \downarrow \\
& \Rightarrow P=\gamma H \tag{1.7}
\end{align*}
$$

Figure 1.2 Pressure developed by a column of fluid.
Observing the equation, it can be concluded that pressure does not depend on the crosssectional area of the liquid column but only on the column height and specific weight of the liquid. Changing the cross-sectional area of the liquid column changes the weight of the liquid by a proportional amount. Hence, $F / A$ (pressure) remains constant.

### 1.4.2 Atmospheric Pressure and Absolute Pressure

Atmospheric pressure is the force per unit area exerted against a surface by the weight of air above that surface in the Earth's atmosphere. In most circumstances, atmospheric pressure is closely approximated by the hydrostatic pressure caused by the weight of air above the measurement point. Low-pressure areas have less atmospheric mass above their location, whereas high-pressure areas have more atmospheric mass above their location. Similarly, as elevation increases, there is less overlying atmospheric mass, such that the pressure decreases with increasing elevation. The standard atmosphere is a unit of pressure that is equal to 101325 Pa or 101.325 kPa . The following units are equivalent, but only to the number of decimal places displayed: 760 mmHg (torr), 29.92 in $\mathrm{Hg}, 14.696 \mathrm{psi}, 1013.25 \mathrm{mbar} / \mathrm{hPa}$. One standard atmosphere is the standard pressure used for pneumatic fluid power (ISO R554) and in the aerospace (ISO 2533) and petroleum (ISO 5024) industries. Absolute pressure is measured relative to a perfect vacuum such as that existing in outer space.

### 1.4.3 Gauge pressure and absolute pressure

Gauge pressure is measured relative to the atmosphere, whereas absolute pressure is measured relative to a perfect vacuum such as that existing in outer space.

A chart showing the difference between gauge and absolute pressure is given in Fig. 1.3. Let us examine the two pressure levels $P_{1}$ and $P_{2}$.

1. Relative to a prefect vacuum, they are:

- $P_{1}=0.7$ bar (absolute), that is, a pressure less than an atmospheric pressure.
- $P_{2}=2$ bar (absolute), that is, a pressure greater than an atmospheric pressure.

2. Relative to atmosphere, they are:

- $P_{1}=-0.3$ bar (gauge, suction or vacuum).
- $P_{2}=1$ bar (gauge).

As can be seen from Fig. 1.3, the following rule can be used in pressure conversion calculations:

$$
\text { Absolute pressure }=\text { Gauge pressure }+ \text { Atmospheric pressure }
$$

It should be noted that vacuum or suction pressures exist in a certain location of fluid power systems (e.g., in the inlet or suction lines of pumps). Therefore, it is important to understand the meaning of pressures below atmospheric pressure. One way to generate a suction pressure is to remove some of the fluid from a closed vessel initially containing fluid at atmospheric pressure.


Figure 1.3 Difference between absolute and gauge pressure.

## Example 1.5

For the fluid power automobile lift system of Fig. 1.4, the hydraulic piston has a $250-\mathrm{mm}$ diameter. How much of oil pressure ( kPa ) is required to lift a 13300 N automobile?


Figure 1.4

Solution: Given force $F=13300 \mathrm{~N}$, diameter $(\mathrm{d})=0.25 \mathrm{~m}$, area $(\mathrm{A})=0.0491 \mathrm{~m}^{2}$. The gauge pressure is
$\operatorname{Pressure}(P)=\frac{F}{A}=\frac{13300}{250}=270.876 \mathrm{kPa}$

## Example 1.6

For the fluid power automobile lift system of Fig. 1.5, the air pressure equals 550 kPagauge. If the hydraulic piston has a diameter of 250 mm , what is the maximum weight of an automobile that can be lifted? The specific gravity of oil is 0.9 . What percentage error in the answer to this problem occurs by ignoring the $1-\mathrm{m}$ head of oil to between the air and interface surface and bottom surface of the piston? Density $(\rho)=900 \mathrm{~kg} / \mathrm{m}^{3}$.


Figure 1.5
Solution: Given pressure of
$\operatorname{air}\left(P_{\text {air }}\right)=550 \mathrm{kPa}$. We know that
$P_{\text {air }}=$ Pressure at piston $\left(P_{\text {piston }}\right)+$ Pressure due to1 moil head

$$
\Rightarrow 550 \times 10^{3}=P_{\text {piston }}+\rho g h
$$

$\Rightarrow P_{\text {piston }}=550 \times 10^{3}-900 \times 9.81 \times 1=541.18 \mathrm{kPa}$
If we ignore 1-m oil head, then $P_{\text {piston }}=550 \mathrm{kPa}$. Then, error in solution is calculated as

$$
\% \text { error }=\frac{550-541.18}{541.18} \times 100=1.63 \%
$$

### 1.5 Compressible and Incompressible Fluids

Fluid power deals with both incompressible and compressible fluids, that is, with oil and air of either constant or variable density. Although there is no such thing in reality as an incompressible fluid, we use this term where the change in density with pressure is so small as to be negligible. This is usually the case with liquids. We may also consider gases to be incompressible when the pressure variation is small compared with the absolute pressure.

Ordinarily, we consider liquids to be incompressible fluids; yet sound waves, which are really pressure waves, travel through them. This provides the evidence of the elasticity of liquids.

The flow of air in a ventilating system is a case where we may treat a gas as incompressible, for the pressure variation is so small that the change in density is of no importance. However, for a gas or steam flowing at a high velocity through a long pipeline, the decrease in pressure may be so high that we cannot ignore the change in density. For an airplane flying at a speed below $100 \mathrm{~m} / \mathrm{s}$, we may consider the air to be of constant density. But as an object moving through air approaches the velocity of sound, which is of the order of $1200 \mathrm{~km} / \mathrm{h}$ depending on temperature, the pressure and density of the air adjacent to the body become materially different from those of the air at some distance away.

### 1.6 Bulk Modulus (Volume Modulus of Elasticity)

Bulk modulus is a measure of the compressibility of a liquid and is required when it is desired to calculate oil volume changes for high pressure and large system volumes such as forging pressures or natural frequencies generally caused by the interaction of fluid compressibility and moving mass. Bulk modulus is analogous to the modulus of elasticity for solids; however, for fluids, it is defined on a volume basis rather than in terms of the familiar one-dimension stress-strain relation. The compressibility (a change in volume due to a change in pressure) of liquid is inversely proportional to its bulk modulus. For liquids, the value of bulk modulus is $1.72 \times 10^{6} \mathrm{kPa}$. The volume modulus of mild steel is about 170000 MPa . Taking a typical value for the volume modulus of cold water to be 2200 MPa , we see that water is about 80 times as compressible as steel. The compressibility of liquids covers a wide range. Mercury, for example, is approximately $8 \%$ as compressible as water, whereas the compressibility of nitric acid is nearly six times that of water.

## Example 1.7

A $500 \mathrm{~cm}^{3}$ sample of oil is to be compressed in a cylinder until its pressure is increased from 1 to 50 atm . If the bulk modulus of oil equals 1750 MPa , find the percentage change in its volume.

Solution: Given volume $(V)=500 \mathrm{~cm}^{3}, P_{1}=1 \mathrm{~atm}, P_{2}=50 \mathrm{~atm}$ and $B=1750 \mathrm{MPa}$. We know that

$$
\begin{aligned}
B= & \frac{-\Delta P}{\Delta V / V} \\
& \Rightarrow \frac{\Delta V}{V}=\frac{-\Delta P}{B}=\frac{-49 \times 1.013 \times 1000000}{1750000000}
\end{aligned}
$$

Reduction in volume $=-0.2836 \%$, which implies that oil is incompressible.

## Example 1.8

A positive displacement pump with a delivery of $1 \mathrm{~L} / \mathrm{min}$ is fed into a pipe with a total volume of 1 L . If the end of the pipe is suddenly blocked, calculate the rise in pressure after 1 s . (The bulk modulus of the fluid being pumped may be taken as 2000 MPa ; neglect any change in the volume of the pipe.)

Solution: Bulk modulus is given as

$$
B=\frac{-\Delta P}{\Delta V / V}
$$

where $\Delta P$ is the change in pressure, $\Delta V$ is the change in volume and $V$ is the original volume. Now

$$
\Delta V=\text { Pump flow in } 1 \mathrm{~s}=\frac{1}{60} \mathrm{~L}
$$

Hence,

$$
\Delta P=B\left(\frac{\Delta V}{V}\right)=\frac{2000}{60}=33.3 \mathrm{MPa}
$$

This rapid rise in pressure illustrates the necessity of having some form of control to limit the rise in pressure in a system should a pump be deadheaded. The control may be built into the pump or may be an external pressure-limiting device such as a relief valve.

## Example 1.9

An 8 L sample of oil is compressed in a cylinder until pressure increases from 0.7 to 2.7 MPa . If the bulk modulus equals 80 MPa , find the change in the volume of oil.

Solution: Given initial volume $V=8 \mathrm{~L}=0.008 \mathrm{~m}^{3}$ and change in pressure $\Delta P=2.7-0.7=2$ MPa. Bulk modulus is given by

$$
B=\frac{-\Delta P}{\Delta V / V}
$$

So, change in volume

$$
\Delta V=-V\left(\frac{\Delta P}{B}\right)=-0.008\left(\frac{2}{80}\right)=-0.002 \mathrm{~m}^{3}
$$

### 1.7 Reynolds Number

It is a dimensionless number referred to as a compressible or incompressible fluid flow. It was postulated by a British engineer Osborne Reynolds. The Reynolds number set criteria by which the fluid flow regime may be distinguished:

$$
R e=\frac{\rho v D}{\mu}
$$

where $\rho$ is the density $\left(\mathrm{kg} / \mathrm{m}^{3}\right), v$ is the velocity of fluid $(\mathrm{m} / \mathrm{s}), D$ is the diameter of the pipe (m) and $\mu$ is the absolute or dynamic viscosity ( Pa s or $\mathrm{ms} / \mathrm{m}^{2}$ ).

### 1.8 Types of Fluid Flow

Based on the range of Reynolds number, the flow of fluid is classified as laminar flow, transition flow and turbulent flow.

1. Laminar flow: In the laminar flow region, the flow is characterized by the smooth motion of the laminae or layers. When there is no macroscopic mixing of adjacent fluid layers for the flow in the laminar regimes, the Reynolds number is less than 2000.
2. Turbulent flow: In the turbulent flow region, the flow is characterized by the random motion of the fluid particles in three dimensions in addition to mean motion. There is considerable macroscopic mixing of adjacent fluid layers and significant velocity fluctuations. For the turbulent flow, the Reynolds number is greater than 4000.
3. Transition flow: In the transition flow region, the flow is in transition between laminar and turbulent flows. The Reynolds number lies between 2000 and 4000.

### 1.9 Ideal Fluid

An ideal fluid is usually defined as a fluid in which there is no friction; it is inviscid (its viscosity is zero). Thus, the internal forces at any section within it are always normal to the section, even during motion. So these forces are purely pressure forces. Although such a fluid does not exist in reality, many fluids approximate frictionless flow at a sufficient distance from solid boundaries, and so we can often conveniently analyze their behaviors by assuming an ideal fluid.

In a real fluid, either liquid or gas, tangential or shearing forces always develop whenever there is a motion relative to body, thus creating fluid friction, because these forces oppose the motion of one particle past another. These friction forces give rise to a fluid property called viscosity.

### 1.10 Viscosity

The viscosity of a fluid is a measure of its resistance to shear or angular deformation. Motor oil, for example, has a high viscosity and resistance to shear; it is cohesive and feels "sticky," whereas gasoline has a low viscosity. The friction forces in a flowing fluid result from the cohesion and momentum which are interchangeable between molecules. The viscosity of typical fluids depends on temperature. Figure 1.7 indicates how the viscosities of typical fluids depend on temperature. As the temperature increases, the viscosities of all liquids decrease, whereas the viscosities of all gases increase. This is because of the force of cohesion, which diminishes with temperature, predominates with liquids, whereas with gases,
the predominating factor is the interchange of molecules between the layers of different velocities. Thus, a rapidly moving gas molecule shifting into a slower moving layer tends to speed up the latter, and a slow-moving molecule entering a faster moving layer tends to slow down the faster moving layer. This molecular interchange sets up a shear or produces a friction force between adjacent layers. At higher temperatures, molecular activity increases, thereby causing the viscosity of gases to increase with temperature.


Figure 1.7Trends in viscosity variation with temperature.
Consider a classic case of two parallel plates (Fig. 1.8), sufficiently large that we can neglect edge conditions, a small distance $Y$ apart, with fluid filling the space in between. The lower plate is stationary whereas the upper one moves parallel to it with a velocity $U$ due to a force $F$ corresponding to some area $A$ of the moving plate.

(a)

(b)

Figure 1.8Velocity profiles across parallel plates: (a) Linear (no bulk flow) and (b) linear (bulk flow to right).

At the boundaries, the particles of a fluid adhere to the walls, so their velocities are zero relative to the walls. This is the so-called no-slip condition that occurs in all viscous fluids. Thus, in Fig. 1.8, the fluid velocities must be $U$ when in contact with the plate at the upper boundary and zero at the lower boundary. We call the form of velocity variation with distance between these two extremes, as depicted in Fig. 1.8, a velocity profile. If the separation distance $Y$ is not too large, the velocity profile is linear, as shown in Fig. 1.8(a). If, in addition, there is a small amount of bulk fluid transport between the plates, which could result from pressure-fed lubrication, for example, the velocity profile becomes the sum of the previous linear profile and a parabolic profile as shown in Fig. 1.8(b); the parabolic additions to the linear profile are zero at the walls and maximum at the centerline. The behavior of the fluid is much as if it is consisted of a series of thin layers, each of which slips a little relative to the next.

For a large class of fluids under the conditions of Fig. 1.8(a), experiments have shown that

$$
F \propto \frac{A U}{Y}
$$

We see from similar triangles that we can replace $U / Y$ by the velocity gradient $d u / d y$. If we now introduce a constant of proportionality $\mu$, we can express the shearing stress $\tau$ between any two thin sheets of fluid by

$$
\begin{equation*}
\tau=\frac{F}{A}=\mu \frac{U}{Y}=\mu \frac{d u}{d y} \tag{1.8}
\end{equation*}
$$

This equation is called Newton's equation of viscosity, since Sir Isaac Newton (1642-1727) first suggested it. Although better known for this formulation of the fundamental laws of motion and gravity for the development of differential calculus, Newton, an English mathematician and natural philosopher, also performed many pioneering studies in fluid mechanics. In transposed form, Eq. (1.8) defines the proportionality constant

$$
\mu=\frac{\tau}{d u / d y}
$$

where $\mu$ is known as the coefficient of viscosity, the dynamic viscosity or simply the viscosity of the fluid. "Absolute viscosity" shall be used to help differentiate it from another viscosity that will be discussed shortly.

In the beginning of this chapter, we noted that the distinction between a solid and a fluid lies in the manner in which each can resist shearing stress. A further distinction among various kinds of fluids and solids will be clarified by referring to Fig. 1.9. In the case of a solid, shear stress depends on the magnitude of the deformation, but Eq. (1.8) shows that in many fluids, shear stress is proportional to the time rate of deformation.


Figure 1.9Velocity profiles - linear (bulk flow to right).
A fluid for which the constant of proportionality does not change with the rate of deformation is called a Newtonian fluid, and this is plotted as a straight line in Fig. 1.9. The slope of this line is the absolute viscosity $\mu$. The ideal fluid with no viscosity falls on the horizontal axis, whereas a true elastic solid is plotted on the vertical axis. A plastic that sustains a certain amount of stress before suffering a plastic flow corresponds to a straight line intersecting the vertical axis at the yield stress. There are certain non-Newtonian fluids in which $\mu$ varies with the rate of deformation. These are relatively uncommon in engineering usage, so we restrict the remainder of this text to the common fluids that under normal conditions obey Newton's equation of viscosity. Note that

$$
\text { Dimension of } \mu=\frac{\text { Dimension of } \tau}{\text { Dimension of }(d u / d y)}=\frac{\left(\mathrm{N} / \mathrm{m}^{2}\right)}{\mathrm{s}^{-1}}=\frac{\mathrm{N} \mathrm{~s}}{\mathrm{~m}^{2}}
$$

A widely used unit of viscosity in the metric system is the poise ( P ) named after Jean Louis Poiseuille (1799-1869). Poiseuille, a French anatomist, was one of the first investigators of viscosity. A poise $=0.10 \mathrm{Ns} / \mathrm{m}^{2}$. A centipoise ( $=0.01 \mathrm{P}$ ) is frequently a more convenient unit. It has a further advantage in that the viscosity of water at $20.2^{\circ} \mathrm{C}$ is 1 cP . Thus, the value of viscosity in centipoise is an indication of the viscosity of a fluid relative to that of water at $20.2^{\circ} \mathrm{C}$.

In many problems involving viscosity, the absolute viscosity is divided by density. This ratio defines the kinematic viscosity, $v$ <COMP: Greek letter nu>, so called because force is not involved, the only dimensions being length and time, as in kinematics. Thus,

$$
\begin{equation*}
v=\frac{\mu}{\rho} \tag{1.9}
\end{equation*}
$$

Usually, the kinematic viscosity $(v)$ is measured in $\mathrm{m}^{2} / \mathrm{s}$. Previously, in the metric system, the common units were $\mathrm{cm}^{2} / \mathrm{s}$, also called the stoke (St), after Sir George Stokes (1819-1903), an English physicist and pioneering investigator of viscosity. Many physicists found centistoke $(\mathrm{cSt})$ a more convenient unit to work with.

A comparison of kinematic viscosities of fluids is shown in Fig. 1.10(a) for low operating pressures. A more detailed characteristic of an ISO 32 mineral oil is shown in Fig. 1.10(b), illustrating the important effect of both temperature and pressure. The data shown in Fig. 1.10 make it clear that experimental testing must specify both pressure and temperature so that related studies may be compared with at least a minimum of confidence. It is common that computer dynamic simulations of hydraulic systems usually assume a mean temperature in a sense that the temperature does not vary significantly during the milliseconds to second of transient behavior. However, it may be necessary to model the effect of pressure on viscosity if large fluctuations in pressure are expected, although its effect may well be of secondary significance.


Figure 1.10Typical kinematic viscosities for a range of fluids. (a) Some fire-resistant fluids. (b) An ISO 32 mineral oil.

## Example 1.11

A 4.5 N force moves a piston inside a cylinder at a velocity of $3 \mathrm{~m} / \mathrm{s}$ as shown in Fig. 1.11. The piston of 10.16 cm diameter is centrally located in the cylinder having an internal diameter of 10.17 cm . An oil film separates the piston from the cylinder. Find the absolute viscosity of the oil.


Figure 1.11
Solution: Given force on the piston $=4.5 \mathrm{~N}$, piston diameter $(d)=10.16 \mathrm{~cm}$, velocity $(v)=3$ $\mathrm{m} / \mathrm{s}$, cylinder diameter $(D)=10.17 \mathrm{~cm}$ and length $(L)=5 \mathrm{~cm}$. Using Newton's law of viscosity we have

$$
F=\mu \frac{d u}{d y}
$$

For small gaps,

$$
\begin{equation*}
F=\mu \frac{A v}{y} \tag{1.10}
\end{equation*}
$$

Now we have

$$
\begin{gathered}
A=\pi D L=\pi \times 10.16 \times 5=159.6 \mathrm{~cm}^{2} \\
v=3 \mathrm{~m} / \mathrm{s}=300 \mathrm{~cm} / \mathrm{s} \\
y=\frac{D-d}{2}=\frac{10.17-10.16}{2}=0.005 \mathrm{~cm}
\end{gathered}
$$

Substituting in Eq. (1.10) we get

$$
4.5=\frac{\mu \times 159.6 \times 300}{0.005}
$$

$$
\Rightarrow \mu=\frac{4.7 \times 10^{-7} \mathrm{~N} \mathrm{~s}}{\mathrm{~cm}^{2}}=\frac{4.7 \times 10^{-7} \mathrm{~N} \mathrm{~s}}{10^{-4} \mathrm{~m}^{2}}=4.7 \times 10^{-3} \mathrm{~N} \mathrm{~s} / \mathrm{m}^{2}
$$

Since $1 \mathrm{~N}=10^{5}$ dynes, we get

$$
\begin{aligned}
\mu & =\frac{4.7 \times 10^{-7} \mathrm{~N} \mathrm{~s}}{\mathrm{~cm}^{2}}=\frac{4.7 \times 10^{-7} \times 10^{5} \text { dynes } \mathrm{s}}{\mathrm{~cm}^{2}} \\
& =0.047 \text { poise }=4.7 \text { centipoise }
\end{aligned}
$$

Now
Absolute viscosity $(\mu)$ in $\mathrm{cP}=$ Specific gravity $(0.89) \times$ Kinetic viscosity $(v)$ in cS
$\Rightarrow v=4.7 / 0.89=5.28 \mathrm{cS}$

## Example 1.12

In Fig. 1.12, oil of the absolute viscosity $\mu$ fills the small gap of thickness $Y$.
(a) Neglecting fluid stress exerted on the circular underside, obtain an expression for the torque $T$ required to rotate the cone at a constant speed $\omega$.
(b) What is the rate of heat generation, in $\mathrm{J} / \mathrm{s}$, if the absolute viscosity of oil is $0.20 \mathrm{Ns} / \mathrm{m}^{2}, \alpha$ $=45^{\circ}, a=45 \mathrm{~mm}, b=60 \mathrm{~mm}, Y=0.2 \mathrm{~mm}$ and the speed of rotation is 90 rpm ?


Figure 1.12

Solution: We have
(a) We have $U=\omega r$. For small gap $Y$ we have

$$
\begin{gathered}
\frac{d u}{d y}=\frac{U}{Y}=\frac{\omega r}{Y} \\
\tau=\frac{\mu d u}{d y}=\frac{\mu \omega r}{Y} \\
d A=2 \pi r d s=\frac{2 \pi r d y}{\cos \alpha} \\
d T=r d F=\frac{2 \pi \mu \omega}{Y \cos \alpha} r^{3} d y \quad(\text { where } r=\tan \alpha) \\
\Rightarrow d T=\frac{\mu \omega r(2 \pi r d y)}{Y(\cos \alpha)} \\
d \mu \omega \tan ^{3} \alpha \\
\cos ^{3} \alpha
\end{gathered} y^{3} d y \quad .
$$

Integrating we get

$$
\begin{gathered}
T=\frac{2 \pi \mu \omega \tan ^{3} \alpha}{Y \cos \alpha} \int_{a}^{a+b} y^{3} d y \\
\Rightarrow T=\frac{2 \pi \mu \omega \tan ^{3} \alpha}{4 Y \cos \alpha}\left[(a+b)^{4}-a^{4}\right]
\end{gathered}
$$

(b) We have

$$
\begin{gathered}
{\left[(a+b)^{4}-a^{4}\right]=(0.105 \mathrm{~m})^{4}-(0.045 \mathrm{~m})^{4}=0.0001175 \mathrm{~m}^{4}} \\
\omega=(90 \mathrm{rev} / \mathrm{min})(2 \pi \mathrm{radians} / \mathrm{rev})(1 \mathrm{~min} / 60 \mathrm{~s})=3 \pi \mathrm{rad} / \mathrm{s}
\end{gathered}
$$

Heat generation rate $=$ Power $=T \omega$

$$
\begin{aligned}
& P=\frac{2 \pi \mu \omega^{2} \tan ^{3} \alpha}{4 Y \cos \alpha}\left[(a+b)^{4}-a^{4}\right] \\
& =\frac{2 \pi\left(0.20 \mathrm{Ns} / \mathrm{m}^{2}\right)\left(3 \pi \mathrm{~s}^{-1}\right)^{2}(1)^{3}\left(0.0001175 \mathrm{~m}^{4}\right)}{4\left(2 \times 10^{-4} \mathrm{~m}\right) \cos 45^{\circ}}
\end{aligned}
$$

$$
=23.2 \mathrm{Nm} / \mathrm{s}=23.2 \mathrm{~J} / \mathrm{s}
$$

### 1.11 Viscosity Index

The viscosity of hydraulic oils decreases with increase in temperature. Hence, the viscosity of the given oil must be represented at a special temperature. The variation of viscosity with respect to temperature is different for different oils.

The viscosity index (VI) is a relative measure of the change in the viscosity of an oil with respect to a change in temperature. An oil having a low VI is one that exhibits a large change in viscosity with a small change in temperature. A high VI oil does not change appreciably with a change in temperature.
The VI of any hydraulic oil can be calculated as follows:

$$
\begin{equation*}
\mathrm{VI}=\frac{L-U}{L-H} \times 100 \tag{1.11}
\end{equation*}
$$

where $L$ is the viscosity in SUS (Saybolt universal viscosity)of a 0 VI oil at $100^{\circ} \mathrm{F}, U$ is the viscosity in SUS of an unknown VI oil at $100^{\circ} \mathrm{F}$ and $H$ is the viscosity in SUS of a 100 VI oil at $100^{\circ} \mathrm{F}$. The VI of an unknown oil is determined from tests. A reference oil of 0 VI and a reference oil of 100 VI are selected. The viscosities of three oils ( $L, U$ and $H$ ) are then measured at $100^{\circ}$ F.This is shown schematically in Fig. 1.13.

A high VI oil is a good all-weather-type oil for use in outdoor machines operating in extreme temperature swings. For a hydraulic system, the oil temperature does not change appreciably; hence, the VI of the oil is not crucial.


Figure 1.13Viscosity index.

## Example 1.13

A sample of oil with viscosity index of 70 is tested with a 0 VI oil and a 100 VI oil whose viscosity values at $100^{\circ} \mathrm{F}$ are 375 and 125 SUS, respectively. What is the viscosity of the sample oil at $100^{\circ} \mathrm{F}$ in units of SUS?

Solution: We know that

$$
\mathrm{VI}=\frac{L-U}{L-H} \times 100
$$

where $L$ is the viscosity of 0 VI oil at $100^{\circ} \mathrm{F}, H$ is the viscosity of 100 VI oil at $100^{\circ} \mathrm{F}$, and $U$ is the required viscosity in SUS So,

$$
\begin{aligned}
& 70=\frac{375-U}{375-125} \times 100 \\
& \Rightarrow U=200 \text { SUS }
\end{aligned}
$$

## Example 1.14

A sample oil is tested with a 0 VI oil and a 100 VI oil whose viscosity values at $38^{\circ} \mathrm{C}$ are 400 and 150 SUS, respectively. If the viscosity of the sample oil at $38^{\circ} \mathrm{C}$ is 200 SSU , what is the viscosity index of the sample oil?

Solution: We know that

$$
\begin{aligned}
\text { VI of the sample oil } & =\frac{L-U}{L-H} \times 100=\frac{400-200}{400-150} \times 100 \\
& =\frac{200}{250} \times 100=80 \mathrm{VI}
\end{aligned}
$$

## Objective Type Questions

## Fill in the Blanks

1. The molecules of a solid are usually $\qquad$ to each other than those of a fluid.
2. The molecules of a gas are much $\qquad$ than those of a liquid.
3. Absolute pressures are measured relative to $\qquad$ such as that existing in outer space.
4. An ideal fluid is usually defined as a fluid in which there is $\qquad$ .
5. The viscosity of a fluid is a measure of its resistance to $\qquad$ .

## State True or False

1. Gauge pressures are measured relative to atmosphere.
2. The bulk modulus is analogous to the modulus of elasticity for solids.
3. At any given temperature, the bulk modulus of water does vary a great deal for a moderate range in pressure.
4.Hydraulic power is especially useful when performing a heavy work.
4. The viscosity of hydraulic oils increases with increase in temperature.
5. Mineral oil has a high bulk modulus compared to phosphate esters and water glycol fluids.

## Review Questions

1. What are the differences between a liquid and a gas?
2. Define the terms specific density, specific weight and specific gravity.
3. Differentiate between absolute and gauge pressures.
4. What is meant by bulk modulus? Give its typical value for liquid and gas.
5. Differentiate between viscosity and viscosity index.
6. Define viscosity index and also suggest an optimum range.
7. Name two undesirable results when using an oil with a viscosity that is too high.
8. Name two undesirable results when using an oil with a viscosity that is too low.
9.What is meant by the term pressure head?
10.Under what conditions is viscosity index important?
11.It is desired to select an oil of a viscosity index suitable for a hydraulic powered front-end loader what will operate a year round. What viscosity index characteristics should be specified? Why?

## Answers

## Fill in the Blanks

1. Closer
2. Farther apart
3. Perfect vacuum
4. No friction
5. Shear or angular deformation

State True or False

1. True
2. True
3. True
4. True
5. False
6.False
